Changing coordinate systems

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The relation between the unit vectors can be written in the symbolic form:

$$\begin{pmatrix} \hat{\mathbf{r}} \\ \hat{\theta} \\ \hat{\varphi} \end{pmatrix} = \mathbf{M} \begin{pmatrix} \hat{\mathbf{x}} \\ \hat{\mathbf{y}} \\ \hat{\mathbf{z}} \end{pmatrix},$$

with $\boldsymbol{\mathsf{M}}$ given by

$$\mathbf{M} = \begin{pmatrix} \sin\theta\cos\phi & \sin\theta\sin\phi & \cos\theta\\ \cos\theta\cos\phi & \cos\theta\sin\phi & -\sin\theta\\ -\sin\phi & \cos\phi & 0 \end{pmatrix}.$$

This way of writing it is not completely legit, since we have "vectors of vectors" but it is a convenient shorthand for eqn. 4.13.

The coordinates $(u_r, u_{\theta}, u_{\varphi})$ and (u_x, u_y, u_z) can be related by plugging the these expressions into $\vec{u} = u_r \hat{r} + u_{\theta} \hat{\theta} + u_{\varphi} \hat{\phi}$:

 $\begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = \begin{pmatrix} \sin\theta\cos\phi & \cos\theta\cos\phi & -\sin\phi \\ \sin\theta\sin\phi & \cos\theta\sin\phi & \cos\phi \\ \cos\theta & -\sin\phi & 0 \end{pmatrix} \begin{pmatrix} u_r \\ u_\theta \\ u_\phi \end{pmatrix}$ And we see that this is the transpose of the previous matrix:

$$\begin{pmatrix} u_{x} \\ u_{y} \\ u_{z} \end{pmatrix} = \mathbf{M}^{\mathsf{T}} \begin{pmatrix} u_{r} \\ u_{\theta} \\ u_{\varphi} \end{pmatrix}$$

Remember the transpose is to swap rows and columns: $M_{ij}^{T} = M_{ji}$.

Now we want the relation to run the other way, we want the spherical coordinates in terms of the cartesian coordinates. Multiply by the inverse $(\mathbf{M}^{\mathsf{T}})^{-1}$:

$$(\mathbf{M}^{T})^{-1} \begin{pmatrix} u_{x} \\ u_{y} \\ u_{z} \end{pmatrix} = (\mathbf{M}^{T})^{-1} \mathbf{M}^{T} \begin{pmatrix} u_{r} \\ u_{\theta} \\ u_{\varphi} \end{pmatrix}$$

matrix times its inverse gives identity:

$$(\mathbf{M}^{\mathsf{T}})^{-1} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = \mathbf{I} \begin{pmatrix} u_r \\ u_\theta \\ u_\varphi \end{pmatrix}$$

which acts like the number 1:

$$(\mathbf{M}^{\mathsf{T}})^{-1} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = \begin{pmatrix} u_r \\ u_\theta \\ u_{\varphi}, \end{pmatrix}$$

So now we have what we want, but what is $(\mathbf{M}^{\mathsf{T}})^{-1}$?

All the columns and rows of **M** are orthogonal (check by dot products). So **M** is an **orthog**-**onal matrix**, which means

$$\mathbf{M}^{-1} = \mathbf{M}^{\mathsf{T}}$$

This can be verified by calculating $\mathbf{M}\mathbf{M}^{\mathsf{T}}$ and $\mathbf{M}^{\mathsf{T}}\mathbf{M}$ (prob 4.2d)

So
$$M^{-1} = M^T$$
, So
 $(M^T)^{-1} = (M^{-1})^{-1} = M$

and so

$$\begin{pmatrix} u_r \\ u_{\theta} \\ u_{\varphi} \end{pmatrix} = \mathbf{M} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix}.$$

Why?

Consider

$$\vec{\mathbf{v}} = v_r \hat{\mathbf{r}} + v_\theta \hat{\theta} + v_\varphi \hat{\phi}$$

in spherical coordinates. If one takes derivatives to see what this is explicitly:

$$\vec{\mathbf{v}} = \dot{\mathbf{r}}\hat{\mathbf{r}} + \mathbf{r}\dot{\theta}\hat{\theta} + \mathbf{r}\sin\theta\dot{\varphi}\hat{\varphi}.$$

This because, for example: $\hat{\mathbf{r}} = \hat{\mathbf{r}}(\theta, \phi)$. The spherical unit vectors are dependent on the coordinates. The basis vectors change as a particle moves.

So this machinery is needed to convert coordinates. The acceleration is even more complex (see Problem 4.3e and equation 4.27.)

Volume integration

$$\iint \int F \, dV = \iint \int F(x, y, z) \, dx dy dz.$$

Volume spanned by three vectors:

volume = det
$$(\vec{a}, \vec{b}, \vec{c}) = \begin{vmatrix} a_x & b_x & c_x \\ a_y & b_y & c_y \\ a_z & b_z & c_z \end{vmatrix}$$
.

Let $\theta \to \theta + d\theta$, then $\vec{r}(r, \theta, \varphi) \to \vec{r}(r, \theta + d\theta, \varphi)$ which gives a change

$$\vec{\mathbf{r}}(\mathbf{r}, \theta + d\theta, \phi) - \vec{\mathbf{r}}(\mathbf{r}, \theta, \phi) = \frac{\partial \vec{\mathbf{r}}}{\partial \theta} d\theta.$$

This will be the infinitesimal vector along the θ direction. Similar reasoning applies to the others. So we get a little volume element:

$$\mathrm{d} \mathbf{V} = \det\left(\frac{\partial \vec{\mathbf{r}}}{\partial r}\,\mathrm{d} r, \frac{\partial \vec{\mathbf{r}}}{\partial \theta}\,\mathrm{d} \theta, \frac{\partial \vec{\mathbf{r}}}{\partial \phi}\,\mathrm{d} \phi\right)$$

The symbolic determinant is the $dV = Jdr d\theta d\phi$, the volume element. It can be written as in problem 4.4b, equation 4.31. The usual shorthand notation for J is:

$$\mathbf{J} = \frac{\partial(\mathbf{x}, \mathbf{y}, z)}{\partial(\mathbf{r}, \theta, \phi)}.$$

If you plug in the expression for \vec{r} and take many partials and work out the determinant for spherical coordinates, you find that.

 $J=r^2\sin\theta,$

So the spherical volume element is

 $dV = r^2 \sin \theta dr d\theta d\varphi.$

This can be done by geometric arguments (as you have probably seen before this) but the Jacobian handles more general coordinate systems.

Cylindrical coordinates

The text shows a clever set of substitutions for deriving equivalent results for cylindrical coordinates, based on the fact that they are the same at the equator of the sphere:

$$\begin{split} \mathbf{r} &= \sqrt{\mathbf{x}^2 + \mathbf{y}^2 + z^2} \rightarrow \sqrt{\mathbf{x}^2 + \mathbf{y}^2}, \\ & \theta \rightarrow \pi/2, \\ & \widehat{\theta} \rightarrow -\widehat{\mathbf{z}}, \\ & \mathbf{r} d\theta \rightarrow - dz. \end{split}$$